Inhomogeneous Static Model in Brans-Dicke Theory

Marcelo Samuel Berman¹

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A static universe with position-dependent rest-energy density, pressure, and scalar field is considered in Brans-Dicke theory. A perfect-gas equation of state is obtained with the solution to the field equations for the Euclidean case with Robertson-Walker metric.

Brans-Dicke (1961) theory is a viable alternative to general relativity when the coupling constant $\omega > 500$. The field equations of this theory are

$$G_{j}^{i} = 8\pi \phi^{-1} T_{j}^{i} + \omega \phi^{-2} (\phi^{i} \phi_{j} - \frac{1}{2} \delta_{j}^{i} \phi^{k} \phi_{k}) + \phi^{-1} (\phi_{j}^{i} - \delta_{j}^{i} \phi_{jk}^{k})$$
(1)

and

$$(3+2\omega)\phi_{ik}^{k} = 8\pi T \tag{2}$$

where G_i^i is the Einstein tensor, T_i^i is the stress-energy tensor, and

$$T \equiv T_k^k \tag{3}$$

$$\phi_i \equiv \frac{\partial \phi}{\partial x^i} \tag{4}$$

For a Robertson-Walker metric for the spatially flat case,

$$ds^{2} = -dt^{2} + a^{2}(t)(dr^{2} + r^{2} d\theta^{2} + r^{2} \sin^{2} \theta d\psi^{2})$$
(5)

the field equations reduce, in the static case, when

$$\dot{a}(t) \equiv 0 \tag{6}$$

¹Department of Physics and Astronomy, University of Alabama, Tuscaloosa, Alabama 35487-0324.

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to the following relations:

$$p\frac{a^{2}}{\phi} = \frac{2}{r}\frac{\phi'}{\phi} - \frac{\omega}{2}\frac{(\phi')^{2}}{\phi^{2}} = \frac{\phi''}{\phi} + \frac{\phi'}{r\phi} + \frac{\omega}{2}\frac{(\phi')^{2}}{\phi^{2}}$$
(7)

$$\rho \frac{a^2}{\phi} = -\frac{\phi''}{\phi} - \frac{2}{r} \frac{\phi'}{\phi} - \frac{\omega}{2} \frac{(\phi')^2}{\phi^2}$$
(8)

$$(3+2\omega)\left(\frac{\phi''}{\phi} + \frac{2\phi'}{r\phi}\right) = (3p-\rho)\frac{a^2}{\phi}$$
(9)

where p and ρ stand for cosmic pressure and rest-energy density, respectively, and T_i^i is for a perfect fluid case.

If we suppose, as a tentative solution,

$$\phi = Ar^n \qquad (A, n \text{ constants}) \tag{10}$$

we find the following solution:

$$\omega = 2/n - 1 \tag{11}$$

$$n = 3 \pm \sqrt{13} \tag{12}$$

$$\rho = \frac{n}{a^2 r^2} \left(4 - \frac{n}{2} \right) \phi \tag{13}$$

Though neither ρ nor p is a constant, they obey a perfect-gas-law equation of state, i.e.,

$$p = \alpha \rho \qquad (\alpha = \text{const})$$
 (14)

This position-dependent solution for ϕ , ρ , and p should be compared with the results obtained in the homogeneous case by Berman *et al.* (1989). For that case, a flat solution required a time-varying ϕ function, given by

$$\phi = -\frac{4\pi\rho_0}{\omega} (t+C)^2 \qquad (C=\text{const})$$
(15)

where

$$\rho = \rho_0 = \text{const} \tag{16}$$

$$p = p_0 = \text{const} \tag{17}$$

Contact with Newton's gravitational constant G can be made by noting that

$$G = \frac{2\omega + 4}{2\omega + 3} \phi^{-1}$$
(18)

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so that we have, in our case,

$$G = \frac{2\omega + 4}{2\omega + 3} A^{-1} r^{-n} \tag{19}$$

Berman *et al.* (1989) commented that the study of static universes is appealing, not only for theoretical reasons, but also because, if there exists an explanation for the observed red-shift other than the expansion of the universe, such study would gain importance. Peratt (1990) has offered insight into such a possible scenario based on plasma physics studies.

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